

Fig. 3 Closed-loop response displaying residual vibration (solid line) compared to open loop (dashed line).

the self-sensing methods depends on proper impedance matching of elements to cancel out the control voltage, this nonlinearity is of crucial importance.

The limit cycle behavior of Fig. 3 can also be explained by non-linearities of the system. Hysteresis, control voltage bleed through, and noise are the three most prevalent causes of suboptimum closed-loop performance for this project. The latter two phenomena were ruled out as causes of the limit cycle behavior after performing numerous simulations. Hysteresis is difficult to simulate but is known to cause limit cycle behavior in control systems. <sup>12</sup> Recall that not only does the actuator exhibit hysteresis, but the self-sensing signal does as well.

#### **Conclusions**

Implementation of the self-sensing control technique using a magnetostrictive actuator is viable. Results of this Note show that significant damping can be added to a structure. However, the effects of nonlinearities inherent in Terfenol-D make the implementation process nontrivial. A good system estimator would have to be used to implement a proper self-sensing control law. Although seemingly complex, additional self-sensing circuitry would eliminate the need for independent, and potentially noncollocated, sensors. Thus, this proof of concept may motivate actuator designers to consider the potential benefits of self-sensing actuation.

#### Acknowledgments

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# Precise Trajectory Tracking Control of Elastic Joint Manipulator

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#### I. Introduction

In the design of high-performance motion systems, the flexibility of the drives and structures, particularly at manipulator arms, presents an obstacle. Some techniques and strategies related to the subject are extensively discussed in Ref. 1.

Similar to the compliance and control of the manipulator arms, a good example of the flexible system is flexible spacecraft, because the location of points at their extremities must be controlled, sometimes to very high precision, by torquing some other point that is separated from the first by sections of flexible structures. In Ref. 3 a survey on the active control technology for large space structures focuses on the development of systematic modeling and design tools for the control of large space structures that has occurred over the past decade.

The objective of this Note is to obtain a transfer function of the servomotor-driven flexible shaft system that includes the natural frequency and damping ratio of the shaft and to propose a trajectory function with which precise trajectory tracking is possible.

The flexible system consists of dc servomotor, harmonic drive, flexible shaft, and manipulator arm. A transfer function relating the actual angular position to the desired angular position of the arm is obtained, including proportional-integral-derivative (PID) type control gains, inductance, flexible system natural frequency, subsystem natural frequency, and the material damping. Transfer function zeros and poles are analyzed in the Laplace domain, and system response to cycloidal input motion is discussed. The close relation between the settling time of the system and the rise time of the trajectory function is presented, and the possibility of the precise tracking of the cycloidal trajectory for the flexible system is shown.

#### II. Formulation

The flexible system considered here consists of a dc servomotor, harmonic drive, flexible shaft, and a manipulator arm. A sketch of the system is shown in Fig. 1a. Also the equivalent system, at which the motor inertia is reflected to the manipulator arm axis, is shown in Fig. 1b. Here, T indicates torque applied by the dc servomotor,  $J_m$  is the motor inertia,  $\theta_m$  is the motor angular displacement, k and c

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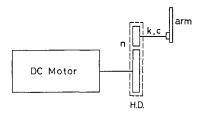


Fig. 1a Dynamic system with flexible shaft.

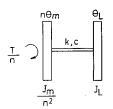


Fig. 1b Equivalent flexible dynamic system.

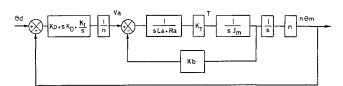


Fig. 2 Closed-loop block diagram of dc servomotor for position con-

are flexible shaft stiffness and damping constant, and  $\theta_L$  and  $J_L$  are arm angle and inertias, respectively. Also, n is the speed reduction ratio of the harmonic drive. For the flexible system shown in Fig. 1b the dynamic equation in matrix form is

$$\begin{bmatrix} J_{m}/n^{2} & 0 \\ 0 & J_{L} \end{bmatrix} \begin{bmatrix} n\ddot{\theta}_{m} \\ \ddot{\theta}_{L} \end{bmatrix} + \begin{bmatrix} c & -c \\ -c & c \end{bmatrix} \begin{bmatrix} n\dot{\theta}_{m} \\ \dot{\theta}_{L} \end{bmatrix} + \begin{bmatrix} k & -k \\ -k & k \end{bmatrix} \begin{bmatrix} n\theta_{m} \\ \theta_{L} \end{bmatrix} = \begin{bmatrix} T/n \\ 0 \end{bmatrix}$$
(1)

The Laplace transform of the equation is

$$\begin{bmatrix} n\theta_m \\ \theta_L \end{bmatrix} = \begin{bmatrix} \left[ \left( J_m / n^2 \right) s^2 + cs + k \right] & -(cs+k) \\ -(cs+k) & \left( J_L s^2 + cs + k \right) \end{bmatrix}^{-1} \begin{bmatrix} T/n \\ 0 \end{bmatrix}$$
(2)

A block diagram of the dc servomotor, which is used for the position control of the manipulator arm, is shown in Fig. 2. Here PID control is assumed. In Fig. 2,  $\theta_d$  is the desired angular displacement;  $K_P$ ,  $K_D$ , and  $K_I$  are proportional, derivative, and integral gains; n is the speed reduction ratio;  $L_a$  and  $R_a$  are motor inductance and resistance;  $K_T$  and  $K_b$  are torque constant and electromotive force constant;  $J_m$  is motor inertia; and  $\theta_m$  is output of the servomotor. From the block diagram, torque T can be obtained in terms of other parameters as

$$\frac{T}{n} = \frac{K_T[K_P + K_D s + (K_I/s)]}{n^2 (sL_a + R_a)} \theta_d$$

$$- \frac{K_T[K_P + K_D s + (K_I/s) + K_b s]}{n^2 (sL_a + R_a)} n\theta_m$$
(3)

When the Eq. (3) is inserted in Eq. (2), the transfer function of the system is obtained as

$$\frac{\theta_L}{\theta_d} = \frac{B(s)}{A(s)} = \frac{b_3 s^3 + b_2 s^2 + b_1 s + b_0}{a_6 s^6 + a_5 s^5 + a_4 s^4 + a_3 s^3 + a_2 s^2 + a_1 s + a_0}$$
(4)

The polynomials B(s) and A(s) are

$$B(s) = \frac{2\zeta_n}{\omega_n} \frac{1}{T_D} s^3 + \left(\frac{2\zeta_n}{\omega_n} \frac{1}{T_P} + \frac{1}{T_D}\right) s^2$$

$$+ \left(\frac{2\zeta_n}{\omega_n} \frac{1}{T_I} + \frac{1}{T_P}\right) s + \frac{1}{T_I}$$

$$A(s) = \frac{L_a}{R_a} \frac{1}{\omega_n^2} s^6 + \left(\frac{L_a}{R_a} \frac{2\zeta_n}{\omega_n} + \frac{1}{\omega_n^2}\right) s^5 + \left[\frac{L_a}{R_a} + \frac{2\zeta_n}{\omega_n} + \frac{1}{\omega_n^2} \left(\frac{1}{T_D} + \frac{1}{T_m}\right)\right] s^4 + \left[1 + \frac{1}{\omega_L^2} \frac{1}{T_P} + \frac{2\zeta_n}{\omega_n} + \frac{1}{T_D} + \frac{1}{T_D}\right] s^3 + \left(\frac{1}{T_D} \frac{1}{T_I} + \frac{2\zeta_n}{\omega_n} \frac{1}{T_P} + \frac{1}{T_D} + \frac{1}{T_m}\right) s^2$$

$$+ \left(\frac{2\zeta_n}{\omega_n} \frac{1}{T_I} + \frac{1}{T_P}\right) s + \frac{1}{T_I}$$

where the mechanical time constant of the system is

$$T_m = \frac{R_a \left( J_m + n^2 J_L \right)}{K_T K_b} \tag{5}$$

 $T_P$ ,  $T_D$ , and  $T_I$  are defined as

$$T_{P} = \frac{R_{a}(J_{m} + n^{2}J_{L})}{K_{T}K_{P}}, \qquad T_{D} = \frac{R_{a}(J_{m} + n^{2}J_{L})}{K_{T}K_{D}}$$

$$T_{I} = \frac{R_{a}(J_{m} + n^{2}J_{L})}{K_{T}K_{I}}$$
(6)

Undamped natural frequency of the flexible system and the subsystem are

$$\omega_n = \sqrt{\frac{J_m + n^2 J_L}{J_m J_L}} k, \qquad \omega_L = \sqrt{\frac{k}{J_L}}$$
 (7)

When the motor side is fixed, the remaining system is defined as the subsystem. Here,  $\zeta_n$  is the material damping ratio. Suppose the shaft is rigid; then  $\omega_n$  and  $\omega_L$  go to infinity and the transfer function reduces to

$$\frac{\theta_L}{\theta_d} = \left(\frac{1}{T_D}s^2 + \frac{1}{T_P}s + \frac{1}{T_I}\right) / \left[\frac{L_a}{R_a}s^4 + s^3 + \left(\frac{1}{T_D} + \frac{1}{T_m}\right)s^2 + \frac{1}{T_P}s + \frac{1}{T_I}\right]$$
(8)

Suppose  $L_a=0$  and torque is not applied to the flexible system (which means  $K_T=0$ ); then the denominator of the transfer function reduces to

$$s^3 \left( s^2 + 2\zeta_n \omega_n s + \omega_n^2 \right) \tag{9}$$

which is the characteristic equation for the two-degree-of-freedom dynamic system.

#### III. Example

To continue the analysis the example given in Table 1 is chosen. All values are taken from Altinay HSR4 model manipulator developed at the KOSGEB facilities of Istanbul Technical University.

 $L_a$  appears in Eq. (4), which is generally neglected with respect to  $R_a$ . Figures 3 and 4 show the root loci of the system with and without the effect of inductance, respectively. For these plots, gains  $K_D$  and  $K_I$  and damping ratio  $\zeta_n$  are set equal to zero. After a certain value of  $K_P$ , 318 for the chosen example, inductance causes instability. When material damping is not considered and gains  $K_D$  and  $K_I$  are set equal to zero, roots start at the poles, which are  $-414.83 \pm 336.27j$ ,  $-14.99 \pm 57.31j$ , 0.0, 0.0, and end at the zeros, which are  $\pm 57.73j$ , 0.0, and infinities. To see the effect of the derivative and integral gains, dominant roots close to the imaginary axis are separately plotted in Fig. 5. Root loci near the imaginary axis not clear in Figs. 3 and 4 can be seen in Fig. 5 as the solid line plot. Results show that roots are more sensitive to  $K_D$  than to  $K_I$ . Increasing values of  $K_D$  rapidly lower the real part of the complex roots, and

Table 1 Characteristics of dc servomotor and manipulator arm

 $R_a = 4.9\Omega$   $L_a = 5.7 \text{ mH}$   $K_T = 0.37 \text{ Nm/A}$   $K_b = 0.37 \text{ Vs/rad}$   $J_m = 0.8 \times 10^{-4} \text{ Nms}^2$   $J_L = 2.4 \text{ Nms}^2$  n = 0.01  $k = 0.8 \times 10^{-4} \text{ Nm/rad}$   $\omega_L = 57.73 \text{ rad/s}$   $\omega_n = 115.47 \text{ rad/s}$ 

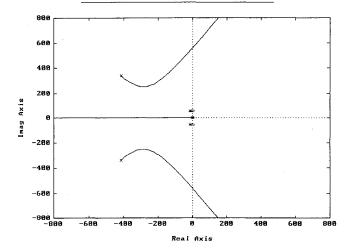


Fig. 3 Root loci when  $L_a$  is not neglected.

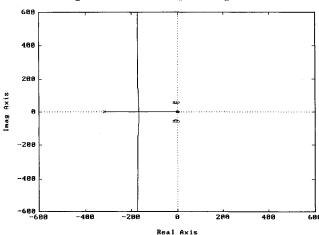


Fig. 4 Root loci when  $L_a$  is neglected.

the system becomes more oscillatory. When the system is assumed rigid, the dominant roots, which characterize the system response, do not appear on the root loci plot; therefore, flexibility should be taken into account in the control of dynamic systems.

The manipulator arm is assumed to track the cycloidal trajectory,<sup>5,6</sup> which is given as

$$\theta_d = \Delta\theta \left(\frac{t}{t_r} - \frac{1}{2\pi}\sin 2\pi \frac{t}{t_r}\right) \tag{10}$$

Here  $\Delta\theta = \theta_2 - \theta_1$ , the angle to be traveled;  $t_r$  is the rise time. Let us take the Laplace transform of Eq. (10)

$$\theta_d(s) = \frac{A}{s_2} - \frac{A}{s^2 + \omega_p^2}, \qquad A = \frac{\Delta\theta\omega_p}{2\pi}, \qquad \omega_p = \frac{2\pi}{t_r}$$
 (11)

Cycloidal motion consists of one ramp and one sinusoidal time function. If the transfer function of the system is type two,<sup>7</sup> then it is expected that the system can follow ramp function without an error. Suppose  $K_P = 20$ ,  $K_D = 0.2$ , and  $K_I = 50$  are selected; then the poles of the transfer function are  $-400.72 \pm 514.91j$ ,  $-6.56 \pm 52.24j$ , -42.38, -2.70. Dominant roots are  $-6.56 \pm 52.24j$ , and

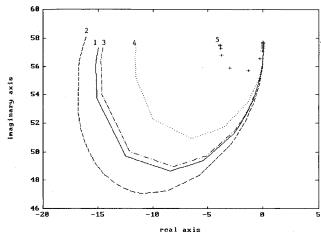


Fig. 5 Root loci of dominant roots: 1)  $K_D = 0$ ,  $K_I = 0$ ; 2)  $K_D = 0$ ,  $K_I = 100$ ; 3)  $K_D = 0.01$ ,  $K_I = 0$ ; 4)  $K_D = 0.1$ ,  $K_I = 0$ ; and 5)  $K_D = 1$ ,  $K_I = 0$ .

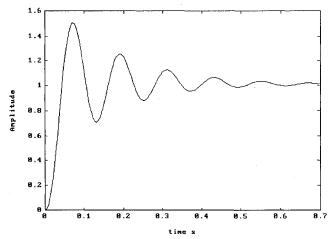


Fig. 6 System response to unit step input.

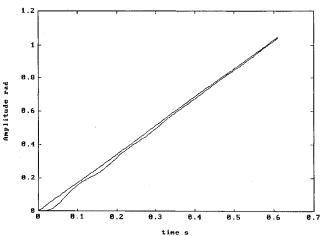


Fig. 7 System response to the ramp part of the cycloidal input.

the system can be considered as a second-order system behaving with these roots. Settling time for the second-order system can be approximated with  $t_s = 4/\sigma$  (Ref. 7). Here  $\sigma$  is the real part of the complex root. For the chosen example,  $t_s = 0.6098$  s. Figure 6 shows the response of the system to a unit step input. As seen from the figure, assumed settling time  $t_s$  is a good approximation for the settling time of the system. Suppose rise time  $t_r$  is chosen as equal to settling time  $t_s$  of the system; then the response of the system to the ramp part of the cycloidal input is given in Fig. 7. As expected, error is almost zero at  $t = t_s = t_r$ . But the same  $t_r$  is not long enough for the sinusoidal part of the cycloidal input to be followed by the system. The response to the sinusoidal input is seen in Fig. 8. It is observed that  $t_r = 4t_s$  is a good approximation for the system to follow the sinusoidal part. For the chosen example, suppose  $t_r = 4t_s = 2.439$  s is assumed; the response of the flexible system to the

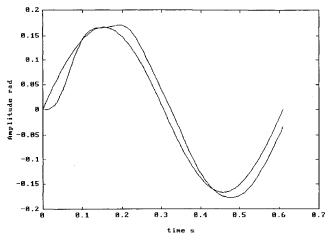


Fig. 8 System response to the sinusoidal part of the cycloidal input.

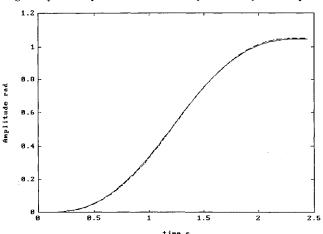


Fig. 9 System response to the cycloidal input function.

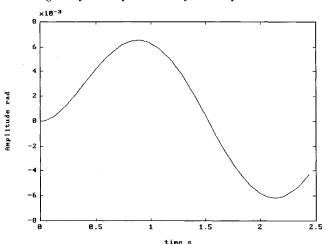


Fig. 10 Tracking error of the flexible system.

cycloidal trajectory and the tracking error are shown in Figs. 9 and 10, respectively. The maximum tracking error is 0.62%. An angular displacement of 1.0472 rad (60 deg) is chosen for the example.

To see the effect of material damping to the system response,  $\zeta_n = 0.05$  is assumed; then for previously chosen values of the gains, the system poles are  $-405.56 \pm 518.68j$ ,  $-7.56 \pm 51.70j$ , -42.25, -2.70. When compared with the poles obtained without material damping, there is little change in system poles and slight increase in the damping of the system, but this does not seriously affect the overall system response.

### IV. Conclusion

In this Note a flexible control system consisting of a dc servomotor, harmonic drive, flexible shaft, and manipulator arm is considered. The transfer function of the system relating the manipulator

arm angle to the desired trajectory angle is derived, including natural frequencies and material damping of the flexible system, as well as PID gains of the control system. It is shown that the dominant roots of the transfer function, which characterize the response of the flexible system, are totally different from that of a rigid system. Material damping is not very effective on the system response. The flexible system is more sensitive to the derivative gain than to the integral gain. For the chosen particular type of trajectory, which is cycloid, the settling time of the system can be related to the rise time of the cycloid, and precise trajectory tracking can be obtained.

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# Exponential Criterion-Based Guidance Law for Acceleration Constrained Missile and Maneuvering Target

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## Introduction

THE linear-quadratic (LQ) and linear-quadratic-Gaussian (LQG) theories have been used to derive modern guidance laws. The structure of such modern guidance laws is a guidance gain multiplied by an estimated zero effort miss. The structure for a high-order missile and a maneuvering target is dealt with in Refs. 1 and 2. The guidance gain depends only on the dynamics of the missile. The estimated zero effort miss is the miss without effort, computed with the estimated states. The estimated states are derived by an appropriate Kalman filter. This structure is because of the separation theorem and certainty equivalence principle.<sup>3</sup>

The linear-exponential and linear-exponential-Gaussian (LEG) problems are generalizations of the LQ and LQG problems. The LEG criterion problem has been solved in Ref. 4. It has been applied to homing missile guidance in Ref. 5. The structure of guidance law based on the LEG criterion for a high-order missile and a high-order maneuvering target is treated in Ref. 6. The effect of the acceleration constraint imposed by the structural or aerodynamic limitations on guidance law, based on the LQ and LQG criteria, for a target whose maneuver is described by a high-order transfer function/shaping filter and a high-order, acceleration constrained missile is described in Refs. 7 and 8.

The structure of a guidance law based on exponential criteria for an arbitrary order and acceleration constrained missile and

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